Scattering Amplitudes in Theories of Compactified Gravity

Dennis Foren (he/they) July 24th, 2020

MICHIGAN STATE UNIVERSITY

UC San Diego

Publications & Dissertation

- NOT Relevant -

"Direct Search Implications for a Custodially-Embedded Composite Top" [Phys. Rev. D 94, 014002] [arXiv:1605.06088]

> "Colorphilic Spin-2 Resonances in the LHC Dijet Channel" [Physica Scripta 93 11 115301, 2018] [arXiv:1706.02502]

- Relevant -

"Scattering Amplitudes of Massive Spin-2 Kaluza-Klein States Grow Only as & (s)"
[Phys. Rev. D 101, 055013] [arXiv:1906.11098]

"Sum Rules for Massive Spin-2 Kaluza-Klein Elastic Scattering Amplitudes" [Phys. Rev. D 100, 115033] [arXiv:1910.06159]

"Massive Spin-2 Scattering Amplitudes in Extra-Dimensional Theories"
[Phys. Rev. D 101, 075013] [arXiv: 2002.12458]

> "Scattering Amplitudes in Theories of Compactified Gravity" [Dissertation]

-

Motivating the Randall-Sundrum 1 (RS1) Model

- Usual 4D Spacetime -



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



- Usual 4D Spacetime -



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



July 24th, 2020

- Usual 4D Spacetime -



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



- Usual 4D Spacetime -



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



July 24th, 2020

- Usual 4D Spacetime -



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



- Usual 4D Spacetime -

- 5D Randall-Sundrum 1 Spacetime -

A five-dimensional (5D) gravity theory with nonfactorizable geometry intro'd by Randall & Sundrum in 1999 to solve the hierarchy problem.



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$



- Usual 4D Spacetime -

- 5D Randall-Sundrum 1 Spacetime -

A five-dimensional (5D) gravity theory with nonfactorizable geometry intro'd by Randall & Sundrum in 1999 to solve the hierarchy problem.



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$

July 24th, 2020

D. Foren – Defense

- Usual 4D Spacetime -

- 5D Randall-Sundrum 1 Spacetime -

A five-dimensional (5D) gravity theory with nonfactorizable geometry intro'd by Randall & Sundrum in 1999 to solve the hierarchy problem.



July 24th, 2020

C(

- Usual 4D Spacetime -

- 5D Randall-Sundrum 1 Spacetime -

A five-dimensional (5D) gravity theory with nonfactorizable geometry intro'd by Randall & Sundrum in 1999 to solve the hierarchy problem.



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$

$$v_{\rm EW} \sim M_{\rm Pl} e^{-\pi k r_c}$$
 $kr_c \sim 12$

D. Foren – Defense

- Usual 4D Spacetime -

- 5D Randall-Sundrum 1 Spacetime -

A five-dimensional (5D) gravity theory with nonfactorizable geometry intro'd by Randall & Sundrum in 1999 to solve the hierarchy problem.



coordinates: $x^{\mu} = (x^0, x^1, x^2, x^3)$

$$v_{\rm EW} \sim M_{\rm Pl} e^{-\pi k r_c}$$

5D-to-4D: 4D Particle Content in RS1



5D RS1 Graviton

5D-to-4D: 4D Particle Content in RS1







5D Diffeomorphism Invariance

July 24th, 2020



Invariance







- New Parameterization of RS1 Model -

- Explicit parameterization of 5D RS1 Lagrangian.
- New term to eliminate all "cosmological constant"-like terms to all orders.
- A 5D-to-4D formula to obtain 4D Effective RS1 Lagrangian.

- Massive Spin-2 KK Mode Scattering -

- Demonstrated O(s) growth in tree-level 2-to-2 massive spin-2 KK mode scattering in RS1 and the 5DOT.
- Derivation & proofs of sum rules (recently generalized to inelastic).
- Simultaneously confirmed numerically.

- Analysis of the 4D Effective RS1 Model -

- Confirmation of 5D strong-coupling scale in the 4D effective model.
- Analysis of KK tower truncation on the 4D effective model accuracy.

(when presenting my original results, I will label w/ corresponding papers info)

From 5D RS1 Lagrangian to 4D Particle Interactions

5D RS1 Metric

$$ds^2 = G_{\mu\nu} \, dx^\mu \, dx^\nu + G_{55} \, dy \, dy$$

$$G_{MN} = \left(\begin{array}{cc} w(x,y) g_{\mu\nu} & 0\\ 0 & -v(x,y)^2 \end{array}\right)$$

$$y = \phi r_c$$

$$y = \phi r_c$$

$$y = \phi r_c$$

- Spin-2 Tower -

- Radion -

$$g_{\mu\nu}(x,y) \equiv \eta_{\mu\nu} + \kappa h_{\mu\nu}(x,y)$$

$$w(x,y) \equiv e^{-2k|y|} e^{-2u(x)}$$

$$u(x,y) \equiv \frac{\kappa \hat{r}(x)}{2\sqrt{6}} e^{+k(2|y| - \pi r_c)}$$

July 24th, 2020

D. Foren – Defense

23 of 95

$$\mathcal{L}_{5D} = \mathcal{L}_{EH} + \mathcal{L}_{CC} + \Delta \mathcal{L}$$

$$\mathcal{L}_{5\mathrm{D}} = \mathcal{L}_{\mathrm{EH}} + \mathcal{L}_{\mathrm{CC}} + \Delta \mathcal{L}$$

- Einstein-Hilbert Lagrangian -

$$\mathcal{L}_{\rm EH} \equiv -\frac{2}{\kappa^2} \sqrt{G} R \qquad \cong -\frac{2}{\kappa^2} \sqrt{G} \tilde{G}^{MN} \left[\Gamma^Q_{MP} \Gamma^P_{NQ} - \Gamma^Q_{MN} \Gamma^P_{PQ} \right]$$

... implies discontinuous curvature at branes.

$$\Gamma_{MN}^{P} \equiv \frac{1}{2} \tilde{G}^{PQ} (\partial_{M} G_{NQ} + \partial_{N} G_{MQ} - \partial_{Q} G_{MN})$$

$$G_{MN} = \begin{pmatrix} w(x, y) g_{\mu\nu} & 0 \\ 0 & -v(x, y)^{2} \end{pmatrix}$$

$$\mathcal{L}_{5\mathrm{D}} = \mathcal{L}_{\mathrm{EH}} + \mathcal{L}_{\mathrm{CC}} + \Delta \mathcal{L}$$

- Cosmological Constant Lagrangian-

$$\mathcal{L}_{\rm CC} \equiv -\frac{2}{\kappa^2} \bigg[-12k^2\sqrt{G} + 6k\sqrt{\overline{G}} \left(\frac{\partial_y^2|y|}{\partial_y^2|y|}\right) \bigg]$$

(bulk CC) + (brane tensions)

$$(\partial_y^2 |y|) = 2 \left[\delta(y) - \delta(y - \pi r_c) \right]$$

$$G_{MN} = \begin{pmatrix} w(x,y) g_{\mu\nu} & 0 \\ 0 & -v(x,y)^2 \end{pmatrix} \quad \overline{G}_{MN} = \begin{pmatrix} w(x,y) g_{\mu\nu} & 0 \\ 0 & 0 \end{pmatrix}$$

July 24th, 2020

D. Foren – Defense

$$\mathcal{L}_{5\mathrm{D}} = \mathcal{L}_{\mathrm{EH}} + \mathcal{L}_{\mathrm{CC}} + \Delta \mathcal{L}$$

- Total Derivative -

[Dissertation] [2002.12458]

$$\Delta \mathcal{L} \equiv \frac{12k}{\kappa^2} \partial_y \left[w^2 \sqrt{-g} \left(\partial_y |y| \right) \right]$$

Ensures...

- CC-like terms cancel
- 2 derivatives per term
- Derivatives act on (different) fields

- <u>A-Type Terms</u> -2 4D derivatives, e.g. $\partial_{\mu}\hat{h}(x,y) \partial_{\nu}\hat{h}(x,y) \hat{h}(x,y)$

€ (E⁺²)

- <u>B-Type Terms</u> -

2 extra-dim. derivatives, e.g. $\partial_y \hat{h}(x, y) \partial_y \hat{h}(x, y) \hat{h}(x, y)$ \overleftarrow{c} (E⁰)

July 24th, 2020

5D RS1 Lagrangian: Weak Field Expansion

$$g_{\mu\nu}(x,y) \equiv \eta_{\mu\nu} + \kappa h_{\mu\nu}(x,y)$$

$$u(x,y) \equiv \frac{\kappa \hat{r}(x)}{2\sqrt{6}} e^{+k(2|y| - \pi r_c)}$$

$$\mathcal{L}_{5\mathrm{D}} = \mathcal{L}_{\mathrm{EH}} + \mathcal{L}_{\mathrm{CC}} + \Delta \mathcal{L}$$

5D RS1 Lagrangian: Weak Field Expansion

$$u(x,y) \equiv \frac{\kappa \hat{r}(x)}{2\sqrt{6}} e^{+k(2|y| - \pi r_c)}$$

$$g_{\mu\nu}(x,y) \equiv \eta_{\mu\nu} + \kappa h_{\mu\nu}(x,y)$$

- Weak Field Expanded (WFE) 5D RS1 Lagrangian -

$$\mathcal{L}_{5D} = \mathcal{L}_{EH} + \mathcal{L}_{CC} + \Delta \mathcal{L} = \mathbf{h} + \mathbf{h}$$

[Dissertation] [2002.12458]

$$= \frac{h}{h} + \frac{r}{r}$$

$$+ \frac{\kappa}{L} \left[\frac{h}{h} + \frac{r}{h} + \frac{r}{h} + \frac{r}{h} + \frac{r}{r} \right]$$

$$+ \frac{\kappa^{2}}{L} \left[\frac{h}{h} + \frac{r}{h} + \frac{r}{h} + \frac{r}{h} + \frac{r}{r} + \frac{r}{r} \right]$$

$$+ \dots$$

5D WFE RS1 Lagrangian

$$\mathcal{L}_{5D} = \frac{h}{h} + \frac{r}{r} + \frac{\kappa}{L} \left[\frac{h}{h} + \frac{r}{h} \right] + \frac{\kappa^2}{L} \left[\frac{h}{h} + \frac{h}{h} \right] + \frac{\kappa^2}{L} \left[\frac{h}{h} + \frac{h}{h} \right] + \dots$$





32 of 95








July 24th, 2020

D. Foren – Defense

36 of 95





$$\hat{h} \hat{h} \hat{h}$$

$$\hat{h}(x,y) \hat{h}(x,y) \hat{h}(x,y) = \frac{1}{(\sqrt{\pi r_c})^3} \sum_{l=0}^{+\infty} \sum_{m=0}^{+\infty} \sum_{n=0}^{+\infty} \begin{bmatrix} \hat{h}^{(l)}(x) \, \hat{h}^{(m)}(x) \, \hat{h}^{(n)}(x) \end{bmatrix} \begin{bmatrix} \psi_l(y) \, \psi_m(y) \, \psi_n(y) \end{bmatrix}$$
4D coordinates extra-dim.

$$\mathcal{L}_{4\mathrm{D}}^{(\mathrm{eff})} = \int_{-\pi r_c}^{+\pi r_c} dy \quad \mathcal{L}_{5\mathrm{D}}$$

$$\hat{h}(x,y) \hat{h}(x,y) \hat{h}(x,y) = \frac{1}{(\sqrt{\pi r_c})^3} \sum_{l=0}^{+\infty} \sum_{m=0}^{+\infty} \sum_{n=0}^{+\infty} \begin{bmatrix} \hat{h}^{(l)}(x) \hat{h}^{(m)}(x) \hat{h}^{(n)}(x) \\ \mathbf{4D \ coordinates} \end{bmatrix} \begin{bmatrix} \psi_l(y) \psi_m(y) \psi_n(y) \\ \mathbf{extra-dim.} \end{bmatrix}$$

$$\mathcal{L}_{4D}^{(\text{eff})} = \int_{-\pi r_c}^{+\pi r_c} dy \quad \mathcal{L}_{5D}$$

$$\frac{1}{(\sqrt{\pi r_c})^3} \sum_{l=0}^{+\infty} \sum_{m=0}^{+\infty} \sum_{n=0}^{+\infty} \begin{bmatrix} \hat{h}^{(l)}(x) \hat{h}^{(m)}(x) \hat{h}^{(n)}(x) \end{bmatrix} \begin{bmatrix} \int_{-\pi r_c}^{+\pi r_c} dy \quad \psi_l(y) \psi_m(y) \psi_n(y) \\ -\pi r_c \end{bmatrix} \begin{bmatrix} \mathbf{1}_{m} \mathbf{1}_{m$$









B-Type = Two Extra-Dimensional Derivatives

 $\partial_y \hat{h}(x,y) \, \partial_y \hat{h}(x,y) \, \hat{h}(x,y)$

€(E⁰)

two extra-dim. derivatives











B-to-A Rules [PROVED]

 $\mu_n \equiv m_n r_c$

[1910.06159] - Elastic Case $b_{n'n'j} = \frac{1}{2} \left[\mu_n^2 - \mu_j^2 \right] a_{nnj}$ $b_{j'n'n} = \frac{1}{2} \mu_j^2 a_{nnj}$ $b_{n'n'nn} = \frac{1}{3} \mu_n^2 a_{nnnn}$

B-to-A Rules [PROVED]

 $\mu_n \equiv m_n r_c$

[1910.06159] - Elastic Case $b_{n'n'j} = \frac{1}{2} \left[\mu_n^2 - \mu_j^2 \right] a_{nnj}$ $b_{j'n'n} = \frac{1}{2} \mu_j^2 a_{nnj}$ $b_{n'n'nn} = \frac{1}{3} \mu_n^2 a_{nnnn}$

[Dissertation] - Inelastic Generalization $b_{l'm'n} = \frac{1}{2} \left[\mu_l^2 + \mu_m^2 - \mu_n^2 \right] a_{lmn}$ $b_{k'l'mn} = \frac{1}{6} \left[2\mu_k^2 + 2\mu_l^2 - \mu_m^2 - \mu_n^2 \right] a_{klmn}$



Elastic Sum Rules [PROVED]

















Elastic Sum Rules [PROVED]

[1910.06159]
$$\mu_n \equiv m_n r$$

[Dissertation]

c



$$c_{nnnn} \equiv \frac{r_c^3}{\pi} \int_{-\pi r_c}^{+\pi r_c} dy \quad e^{-6k|y|} (\partial_y \psi_n)^4$$

[1910.06

Elastic Sum Rules [PROVED]

[1910.06159]
$$\mu_n \equiv m_n r_c$$

[Dissertation]



$$c_{nnnn} \equiv \frac{r_c^3}{\pi} \int_{-\pi r_c}^{+\pi r_c} dy \quad e^{-6k|y|} (\partial_y \psi_n)^4$$

[1910.06]

$$\sum_{j} \left[\mu_{j}^{2} - 5\mu_{n}^{2} \right] \mu_{j}^{4} a_{jnn}^{2} = -\frac{16}{3} \mu_{n}^{6} a_{nnnn}$$

Inelastic Sum Rules [PROVED]

$$\begin{split} \sum_{j} a_{jkl} a_{jmn} &= a_{klmn} \\ \sum_{j} \mu_{j}^{2} a_{jkl} a_{jmn} &= \frac{1}{3} \vec{\mu}^{2} \underline{a_{klmn}} \\ \sum_{j} \mu_{j}^{4} a_{jkl} a_{jmn} &= 4 \underline{c_{klmn}} + \left[\frac{1}{3} (\vec{\mu}^{2})^{2} - (\mu_{k}^{2} + \mu_{l}^{2}) (\mu_{m}^{2} + \mu_{n}^{2}) \right] \underline{a_{klmn}} \\ \sum_{j} \mu_{j}^{6} a_{jkl} a_{jmn} &= 5 \vec{\mu}^{2} \underline{c_{klmn}} - \frac{1}{9} \left[6 (\mu_{k}^{4} + \mu_{l}^{4}) (\mu_{m}^{2} + \mu_{n}^{2}) + 6 (\mu_{k}^{2} + \mu_{l}^{2}) (\mu_{m}^{4} + \mu_{n}^{4}) \right. \\ &\left. - 4 (\mu_{k}^{2} + \mu_{l}^{2})^{3} - 4 (\mu_{m}^{2} + \mu_{n}^{2})^{3} + (\mu_{k}^{6} + \mu_{l}^{6} + \mu_{m}^{6} + \mu_{n}^{6}) \right] \underline{a_{klmn}} \end{split}$$

[Dissertation]

 $\mu_n \equiv m_n r_c$

Matrix Elements

Matrix Element: General Considerations



External States: Massive Spin-2 KK Modes



The *n*th massive spin-2 KK mode $\hat{h}^{(n)}_{\mu
u}$ has <u>5 available helicities</u>

$$\lambda \in \{ -2, -1, 0, +1, +2 \}$$

External States: Massive Spin-2 KK Modes



External States: Massive Spin-2 KK Modes

The *nth* massive spin-2 KK mode $\hat{h}^{(n)}_{\mu
u}$ has <u>5 available helicities</u>

 $\lambda \in \{ -2, -1, 0, +1, +2 \}$



Matrix Element: Elastic Helicity-Zero Process



Contact: Elastic Helicity-Zero Process



Radion Mediated: s - channel



Radion Mediated: t - and u - channel



Spin-2 Mediated: s - channel



Spin-2 Mediated: s - channel



Spin-2 Mediated: s - channel



Matrix Element: Total Matrix Element

$$\mathcal{M}_{\text{full}} = \prod_{n \neq \infty}^{n} \sum_{n \neq \infty}^{\infty} \prod_{n \neq \infty}^{n} = \mathcal{M}_{c} + \mathcal{M}_{r} + \sum_{j=0}^{+\infty} \mathcal{M}_{j}$$





Helicity-Zero Elastic Process: Cancellations [1906.11098] [1910.06159] [2002.12458]












Helicity-Zero Elastic Process: Cancellations [1906.11098]



Helicity-Zero Elastic Process: Cancellations [1906.11098]



Helicity-Zero Elastic Process: Cancellations [1906.11098]



∂ (s⁵) Sum Rule

🔁 (s⁴) Sum Rule

$$\sum_{j} a_{jnn}^2 = a_{nnnn}$$
$$\sum_{j} \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$$

🔁 (s³) Sum Rule

$$\sum_{j=0}^{+\infty} \mu_j^4 a_{nnj}^2 = \frac{16}{15} \mu_n^4 a_{nnnn} + \frac{4}{5} \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$$

🔁 (s²) Sum Rule

$$\sum_{j=0}^{+\infty} \left[\mu_j^2 - \frac{5}{2} \mu_n^2 \right] \mu_j^4 a_{nnj}^2 = -\frac{8}{3} \mu_n^6 a_{nnnn} + 2\mu_n^2 \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$$

July 24th, 2020

[1910.06159]

🔁 (s⁵) Sum Rule **c** (s⁴) Sum Rule $\sum_{j} a_{jnn}^2 = a_{nnnn}$ $\sum_{i} \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$ 🔁 (s³) Sum Rule $\sum_{j=0}^{+\infty} \mu_j^4 a_{nnj}^2 = \frac{16}{15} \mu_n^4 a_{nnnn} + \frac{4}{\xi} \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$ **C** (s²) Sum Rule $\sum_{i=0}^{+\infty} \left[\mu_j^2 - \frac{5}{2} \mu_n^2 \right] \mu_j^4 a_{nnj}^2 = -\frac{8}{3} \mu_n^6 a_{nnn} + 2\mu_i^4 \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$

July 24th, 2020

D. Foren – Defense

[1910.06159]

∂ (s⁵) Sum Rule

🔁 (s⁴) Sum Rule

$$\sum_{j} a_{jnn}^2 = a_{nnnn} \qquad \qquad \sum_{j} \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$$

c (s³) - **c** (s²) Sum Rule

$$\sum_{j} \left[\mu_{j}^{2} - 5\mu_{n}^{2} \right] \mu_{j}^{4} a_{jnn}^{2} = -\frac{16}{3} \mu_{n}^{6} a_{nnnn}$$

🔁 (s²) Sum Rule

$$\sum_{j=0}^{+\infty} \left[\mu_j^2 - \frac{5}{2} \mu_n^2 \right] \mu_j^4 a_{nnj}^2 = -\frac{8}{3} \mu_n^6 a_{nnnn} + 2\mu_n^2 \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$$

July 24th, 2020

[1910.06159]

🔁 (s⁵) Sum Rule

🔁 (s⁴) Sum Rule

$$\sum_{j} a_{jnn}^2 = a_{nnnn}$$

$$\sum_{j} \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$$

PROVED

[1910.06159]

[2002.12458]

ð (s³) - **ð** (s²) Sum Rule

$$\sum_{j} \left[\mu_{j}^{2} - 5\mu_{n}^{2} \right] \mu_{j}^{4} a_{jnn}^{2} = -\frac{16}{3} \mu_{n}^{6} a_{nnnn}$$

PROVED

🔁 (s²) Sum Rule

$$\sum_{j=0}^{+\infty} \left[\mu_j^2 - \frac{5}{2} \mu_n^2 \right] \mu_j^4 a_{nnj}^2 = -\frac{8}{3} \mu_n^6 a_{nnnn} + 2\mu_n^2 \left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2 \right]$$

July 24th, 2020

Helicity-Zero Elastic Process: Sum Rules[1910.06159] $\mathfrak{O}(s^5)$ Sum Rule $\mathfrak{O}(s^4)$ Sum Rule $\mathfrak{O}(s^5)$ Sum Rule $\mathfrak{O}(s^4)$ Sum Rule**PROVED** $\sum_j a_{jnn}^2 = a_{nnnn}$ $\sum_j \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$ **PROVED**

🔁 (s³) - 🔁 (s²) Sum Rule

$$\sum_{j} \left[\mu_{j}^{2} - 5\mu_{n}^{2} \right] \mu_{j}^{4} a_{jnn}^{2} = -\frac{16}{3} \mu_{n}^{6} a_{nnnn}$$

PROVED

Simplified *H* (s²) Sum Rule

$$3\left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2\right] = 15c_{nnnn} + \mu_n^4 a_{nnnn}$$

D. Foren – Defense

Helicity-Zero Elastic Process: Sum Rules[1910.06159]
[2002.12458]
[Dissertation] \mathfrak{O} (s⁵) Sum Rule \mathfrak{O} (s⁴) Sum Rule[Dissertation]**PROVED** $\sum_{j} a_{jnn}^2 = a_{nnnn}$ $\sum_{j} \mu_j^2 a_{jnn}^2 = \frac{4}{3} \mu_n^2 a_{nnnn}$ **PROVED**

🔁 (s³) - 🔁 (s²) Sum Rule

$$\sum_{j} \left[\mu_{j}^{2} - 5\mu_{n}^{2} \right] \mu_{j}^{4} a_{jnn}^{2} = -\frac{16}{3} \mu_{n}^{6} a_{nnnn}$$

PROVED

Simplified *H* (s²) Sum Rule

$$3\left[9b_{n'n'r}^2 - \mu_n^4 a_{nn0}^2\right] = 15c_{nnnn} + \mu_n^4 a_{nnnn}$$



July 24th, 2020

D. Foren – Defense

Helicity-Zero Elastic Process: O(s) Growth

[1906.11098] [1910.06159] [2002.12458]

$$\mathcal{H}_{c} + \sum_{i} \mathcal{H}_{c} + \mathcal{H}_{c} = \mathcal{H}_{c}$$
$$\mathcal{M}_{c} - \mathcal{M}_{j} - \mathcal{M}_{r} - \mathcal{M}_{full}$$
$$\mathcal{O}_{c} \mathbf{s} + \mathcal{O}_{j} \mathbf{s} + \mathcal{O}_{r} \mathbf{s} = \mathcal{M}_{full}$$

$$\overline{\mathcal{M}}^{(1)} = \frac{\kappa^2 \left[7 + \cos(2\theta)\right]^2 \csc^2 \theta}{2304 \, \pi r_c} \left\{ \sum_j \frac{m_j^8}{m_n^8} a_{nnj}^2 + \frac{28}{15} a_{nnnn} - \frac{48}{5} \left[\frac{9 \, b_{n'n'r}^2}{(m_n r_c)^4} - a_{nn0}^2 \right] \right\}$$



July 24th, 2020

 $\mathcal{M}^{[N]}$

 \equiv

[2002.12458]



July 24th, 2020

 $\mathcal{M}^{[N]}$

Strong Coupling Scale Is it consistent with expectations?

$$\Lambda_{\pi} \equiv M_{\rm Pl} \, e^{-\pi k r_C}$$

Strong Coupling Scale



July 24th, 2020

D. Foren – Defense

[2002.12458]







Finite Truncation at Low Energy How many states should I include? [2002.12458]

Truncation at Low Energies: $E = 10 m_1$



July 24th, 2020

D. Foren – Defense

[2002.12458]

92 of 95

1

X

Conclusion

Future Work

- Bulk/Brane Matter -

- Radion Stabilization -









Radion can be **stabilized** through mixing:



July 24th, 2020

D. Foren – Defense

Conclusion

- Contact Information -

- Further Reading -

"Scattering Amplitudes of Massive Spin-2 Kaluza-Klein States Grow Only as *C*(s)"
[Phys. Rev. D 101, 055013] [arXiv:1906.11098]

"Sum Rules for Massive Spin-2 Kaluza-Klein Elastic Scattering Amplitudes" [Phys. Rev. D 100, 115033] [arXiv:1910.06159]

"Massive Spin-2 Scattering Amplitudes in Extra-Dimensional Theories" [Phys. Rev. D 101, 075013] [arXiv: 2002.12458]

> "Scattering Amplitudes in Theories of Compactified Gravity" [Dissertation]



Thank you for your time and attention!

July 24th, 2020